Bogoliubov Laboratory of Theoretical Physics
JOINT INSTITUTE FOR NUCLEAR RESEARCH
141980 Dubna (Moscow region), Russia

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A.K.Motovilov[1]

REMOVAL OF THE ENERGY DEPENDENCE FROM THE RESOLVENT-LIKE ENERGY-DEPENDENT INTERACTIONS[2]

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[1]E–mail: MOTOVILV@THSUN1.JINR.DUBNA.SU.
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Removal of the Energy Dependence from the Resolvent-like Energy-Dependent Interactions

The spectral problem $(A+V(z))\psi = z\psi$ is considered with A, a self-adjoint Hamiltonian of sufficiently arbitrary nature. The perturbation $V(z)$ is assumed to depend on the energy z as resolvent of another self-adjoint operator $A' : V(z) = -B(A'-z)^{-1}B^*$. It is supposed that operator B has a finite Hilbert-Schmidt norm and spectra of operators A and A' are separated. The conditions are formulated when the perturbation $V(z)$ may be replaced with an energy-independent "potential" W such that the Hamiltonian $H = A + W$ has the same spectrum (more exactly a part of spectrum) and the same eigenfunctions as the initial spectral problem. The orthogonality and expansion theorems are proved for eigenfunction systems of the Hamiltonian $H = A + W$. Scattering theory is developed for H in the case when operator A has continuous spectrum. Applications of the results obtained to few-body problems are discussed.

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1. INTRODUCTION

Perturbations, depending on the spectral parameter (usually energy of system) arise in a lot of physical problems (see papers [1]—[16] and Refs. therein). In particular, such are the interaction potentials between clusters formed by quantum particles [1]—[6].

The perturbations of this type appear typically [1]—[4], [11]—[16] as a result of dividing the Hilbert space $\mathcal{H}$ of physical system in two subspaces, $\mathcal{H} = \mathcal{H}_1 \oplus \mathcal{H}_2$. The first one, say $\mathcal{H}_1$, is interpreted as a space of "external" (for example, hadronic) degrees of freedom. The second one, $\mathcal{H}_2$, is associated with an "internal" (for example, quark) structure of the system. The Hamiltonian $\mathbf{H}$ of the system looks as a matrix,

$$\mathbf{H} = \begin{bmatrix} A_1 & B_{12} \\ B_{21} & A_2 \end{bmatrix} \tag{1.1}$$

with A_1, A_2, the channel Hamiltonians (self-adjoint operators) and B_{12}, $B_{21} = B_{12}^*$, the coupling operators. Reducing the spectral problem $\mathbf{H}U = zU$, $U = \{u_1, u_2\}$ to the channel α only one gets the spectral problem

$$[A_\alpha + V_\alpha(z)]u_\alpha = zu_\alpha, \quad \alpha = 1, 2, \tag{1.2}$$

where the perturbation

$$V_\alpha(z) = -B_{\alpha\beta}(A_\beta - z)^{-1}B_{\beta\alpha}, \quad \beta \neq \alpha, \tag{1.3}$$

depends on the spectral parameter z as the resolvent $(A_\beta - z)^{-1}$ of the Hamiltonian A_β. In more complicated cases $V_\alpha(z)$ can include also linear terms in respect with z. Other types of dependency of the potentials $V_\alpha(z)$ on the spectral parameter z give, in a general way, the spectral problems (1.2) with a complex spectrum.

The present paper is a continuation of the author's works [17]—[19] devoted to a study of the possibility to "remove" the energy dependence from perturbations of the type (1.3). Namely, in [17]—[19] we construct such new potentials W_α that spectrum of the Hamiltonian $H_\alpha = A_\alpha + W_\alpha$ is a part of the spectrum of the problem (1.2). At the same time, the respective eigenvectors of H_α become also those for (1.2). Hamiltonians H_α are found as solutions of the non-linear operator equations

$$H_\alpha = A_\alpha + V_\alpha(H_\alpha) \tag{1.4}$$

first appeared in the paper [9] by M.A.Braun in connection with consideration of the quasipotential equation. The operator-value function $V_\alpha(Y)$ of the operator variable Y, $Y : \mathcal{H}_\alpha \to \mathcal{H}_\alpha$, is defined by us in such a way (see Sec. 3) that eigenvectors ψ of Y, $Y\psi = z\psi$, become automatically those for $V_\alpha(Y)$ and $V_\alpha(Y)\psi = V_\alpha(z)\psi$.

In Ref. [17], the case is considered in details when one of the operators A_α is the Schrödinger operator in $L_2(\mathbf{R}^n)$ and another one has a discrete spectrum only. The reports [18], [19] announce the results concerning the equations (1.4) and properties of their solutions H_α in a rather more general situation when the Hamiltonian $\mathbf{H}$ may be rewritten in terms of a two-channel variant of the Friedrichs model investigated by O.A.Ladyzhenskaya and L.D.Faddeev in Refs. [20], [21]. In Ref. [16] the method [17]—[19] is used to construct an effective cluster Hamiltonian for atoms adsorbed by the metal surface.

In the present paper, we specify the assertions from [18], [19] and give proofs for them. Also, we pay attention to an important circumstance disclosing a nature of solutions of the basic equations (1.4). Thing is that the potentials $W_\alpha = V_\alpha(H_\alpha)$ may be presented in the form $W_\alpha = B_{\alpha\beta}Q_{\beta\alpha}$ where the operators $Q_{\beta\alpha}$ satisfy the equations (3.13) (see Sec. 3).

Exactly the same equations arise in the method of construction of invariant subspaces for self-adjoint operators developed by V.A.Malyshev and R.A.Minlos in Refs. [22], [23]. It follows from the results of [22], [23] that operators H_α, $\alpha = 1, 2$, determine in fact, parts of the two-channel Hamiltonian $\mathbf{H}$ acting in corresponding invariant subspaces (see Theorem 2 and comments to it).

Recently, the author came to know about the work[2] "Spectral properties of a class of rational operator-value functions" by V.M.Adamyan and H.Langer studying the operator-value functions written in our notation as $F_\alpha(z) = z - A_\alpha \pm B_{\alpha\beta}(A_\beta - z)^{-1} B_{\beta\alpha}$. In particular Adamyan and Langer show in this work that a subset of eigenvectors of F_α can be chosen to form a Riesz basis in $\mathcal{H}_\alpha$. There is a certain intersection of their results and ours from Refs. [17]—[19]. However the methods are different.

The paper is organized as follows.

In Sec. 2 we describe the Hamiltonian $\mathbf{H}$ as a two-channel variant of the Friedrichs model [20], [21]. We suppose that both operators A_α, $\alpha = 1, 2$, may have continuous spectrum. When properties of objects connected with this spectrum (wave operators and scattering matrices) are considered in following sections, the coupling operators $B_{\alpha\beta}$ in (1.1) are assumed to be integral ones with kernels $B_{\alpha\beta}(\lambda, \mu)$, the Hölder functions in both variables λ, μ.

In Sec. 3 the equations (1.4) are studied. As in Refs. [22], [23] we suppose that spectra $\sigma(A_1)$ and $\sigma(A_2)$ of the operators A_1 and A_2 are separated, $\mathrm{dist}\{\sigma(A_1), \sigma(A_2)\} > 0$. Existence of solutions of Eqs. (1.4) is established only in the case when the Hilbert-Schmidt norm $\|B_{\alpha\beta}\|_2$ of the coupling operators satisfies the condition $\|B_{\alpha\beta}\|_2 < \frac{1}{2}\mathrm{dist}\{\sigma(A_1), \sigma(A_2)\}$.

In Sec. 4 the eigenfunctions systems of the operators H_α are studied and theorems of their orthogonality and completeness are proved. We show here in particular that spectrum of the Hamiltonian $\mathbf{H}$ is distributed between the solutions $H_1 = A_1 + B_{12}Q_{21}$ and $H_2 = A_2 + B_{21}Q_{12}$, $Q_{21} = -Q_{12}^*$, of the basic equations (1.4) in such a way that H_1 and H_2 have not "common" eigenfunctions $U = \{u_1, u_2\}$ of $\mathbf{H}$: simultaneously, component u_1 can not be eigenfunction for H_1, and component u_2, for H_2.

In Sec. 5 we introduce new inner products in the Hilbert spaces $\mathcal{H}_\alpha$, $\alpha = 1, 2$, making the Hamiltonians H_α self-adjoint.

In Sec. 6 we give a non-stationary formulation of the scattering problem for a system described by the Hamiltonians H_α constructed. We show that this formulation is correct and scattering operator is exactly the same as in initial spectral problem.

At last, in Sec. 7 we discuss the questions concerning a use of two-body energy-dependent potentials in few-body problems.

2. INITIAL SPECTRAL PROBLEM AND TWO-CHANNEL HAMILTONIAN

Let A_1 and A_2 be self-adjoint operators acting, respectively, in "external", $\mathcal{H}_1$, and "internal", $\mathcal{H}_2$, Hilbert spaces. We study the spectral problem (1.2) with perturbation $V_\alpha(z)$ given by (1.3). We suppose that $B_{\alpha\beta} \in \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\beta)$ where $\mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\beta)$ is the Banach space of bounded linear operators acting from $\mathcal{H}_\alpha$ to $\mathcal{H}_\beta$.

Note that method developed in the present paper works also in the case of more

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general perturbations[3] $V_\alpha(z) = -\mathcal{R}_\alpha(z)$ containing linear terms,

$$\mathcal{R}_\alpha(z) = N_\alpha z + B_{\alpha\beta}(A_\beta - N_\beta z - z)^{-1} B_{\beta\alpha} \tag{2.1}$$

with N_α, self-adjoint bounded operator in $\mathcal{H}_\alpha$ such that $N_\alpha \geq (\delta - 1) I_\alpha$ where $\delta > 0$ and I_α is the identity operator in $\mathcal{H}_\alpha$. Thing is that the equation (1.2) with $V_\alpha(z) = -\mathcal{R}_\alpha(z)$ can be easily rewritten in the form (1.2),(1.3). To do this, one has only to make the replacements $u_\alpha \to u'_\alpha = (I_\alpha + N_\alpha)^{1/2} u_\alpha$, $A_\alpha \to A'_\alpha = (I_\alpha + N_\alpha)^{-1/2} A_\alpha (I_\alpha + N_\alpha)^{-1/2}$ and $B_{\alpha\beta} \to B'_{\alpha\beta} = (I_\alpha + N_\alpha)^{-1/2} B_{\alpha\beta} (I_\beta + N_\beta)^{-1/2}$. Therefore we shall consider further only the initial spectral problem (1.2),(1.3).

We shall assume that operators A_α, $\alpha = 1, 2$, may have continuous spectra σ^c_α. To deal with these spectra we accept below some presuppositions with respect to A_α restricting us to the case of a two-channel variant of the Friedrichs model [20], [21]. Note that these presuppositions are not necessary for a part of statements (Lemma 1, Theorems 1 — 3 and 5) which stay correct also in general case.

The presuppositions are following.

At first, we assume that Hamiltonian H is defined in that representation where operators $A_\alpha, \alpha = 1, 2$, are diagonal. We suppose that continuous spectra σ^c_α of the operators $A_\alpha, \alpha = 1, 2$, are absolutely continuous and consist of a finite number of finite (and may be one or two infinite) intervals $(a^{(j)}_\alpha, b^{(j)}_\alpha)$, $-\infty \leq a^{(j)}_\alpha < b^{(j)}_\alpha \leq +\infty$, $j = 1, 2, \ldots, n_\alpha$, $n_\alpha < \infty$. At second, we suppose that discrete spectra σ^d_α of the operators A_α, $\alpha = 1, 2$, do not intersect with σ^c_α, $\sigma^d_\alpha \bigcap \sigma^c_\alpha = \emptyset$, and consist of a finite number of points with finite multiplicity. In this case the space $\mathcal{H}_\alpha$ may be present as the direct integral [25]

$$\mathcal{H}_\alpha = \sum\!\!\!\!\!\!\int_{\lambda \in \sigma_\alpha} \oplus \mathcal{G}_\alpha(\lambda) d\lambda \equiv \bigoplus_{\lambda \in \sigma^d_\alpha} \mathcal{G}_\alpha(\lambda) \oplus \int_{\lambda \in \sigma^c_\alpha} \oplus \mathcal{G}_\alpha(\lambda) d\lambda, \quad \sigma_\alpha = \sigma^c_\alpha \bigcup \sigma^d_\alpha \subset \mathbf{R}. \tag{2.2}$$

The space $\mathcal{H}_\alpha$ consists of the measurable functions f_α which are defined on σ_α and have the values $f_\alpha(\lambda)$ from corresponding Hilbert spaces $\mathcal{G}_\alpha(\lambda)$. By $\langle \cdot\,, \cdot \rangle$ we denote the inner product in $\mathcal{H}_\alpha$,

$$\langle f_\alpha, g_\alpha \rangle = \sum\!\!\!\!\!\!\int_{\lambda \in \sigma_\alpha} (f_\alpha(\lambda), g_\alpha(\lambda)) \equiv \sum_{\lambda \in \sigma^d_\alpha} (f_\alpha(\lambda), g_\alpha(\lambda)) + \int_{\lambda \in \sigma^c_\alpha} d\lambda (f_\alpha(\lambda), g_\alpha(\lambda)),$$

where $(\cdot\,, \cdot)$ stands for inner product in $\mathcal{G}_\alpha(\lambda)$. By $|\cdot|$ we denote norm of vectors and operators in $\mathcal{G}_\alpha(\lambda)$ and by $\|\cdot\|$, the norm in $\mathcal{H}_\alpha$. Operator A_α acts in $\mathcal{H}_\alpha$ as the independent variable multiplication operator,

$$(A_\alpha f_\alpha)(\lambda) = \lambda \cdot f_\alpha(\lambda), \quad \alpha = 1, 2. \tag{2.3}$$

It's domain $\mathcal{D}(A_\alpha)$ consists of those functions $f_\alpha \in \mathcal{H}_\alpha$ which satisfy the condition $\sum\!\!\!\!\!\!\int_{\lambda \in \sigma_\alpha} \lambda^2\, |f_\alpha(\lambda)|^2 < \infty$. For the sake of simplicity we assume that $\mathcal{G}_\alpha(\lambda)$ does not depend on $\lambda \in \sigma^c_\alpha$, i.e. $\mathcal{G}_\alpha(\lambda) \equiv \mathcal{G}^c_\alpha$ for each $\lambda \in \sigma^c_\alpha$. Hence, $\int_{\sigma^c_\alpha} \oplus \mathcal{G}_\alpha(\lambda) d\lambda = L_2(\sigma^c_\alpha, \mathcal{G}^c_\alpha) \equiv \mathcal{H}^c_\alpha$. By $E_\alpha(d\lambda)$ we denote a spectral measure [25] of the operator A_α, $A_\alpha = \int_{\sigma_\alpha} \lambda E_\alpha(d\lambda)$. In the diagonal representation considered, the spectral projector E_α acts on $f \in \mathcal{H}_\alpha$ as

$$(E_\alpha(\Delta) f)(\lambda) = \chi_\Delta(\lambda) f(\lambda) \tag{2.4}$$

[3]Remember that if $N_\alpha \geq 0$ then Eq. (2.1) gives a general form of R-function on $\mathcal{H}_\alpha$, i.e. an analytic at $\operatorname{Im} z \neq 0$ $\mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha)$–value function with positive imaginary part for $z : \operatorname{Im} z > 0$ (see paper [24] and Refs. therein).

for any Borelian set $\Delta \subset \sigma_\alpha$. Here, χ_Δ is a characteristic function of Δ, $\chi_\Delta(\lambda) = 1$ if $\lambda \in \Delta$, and $\chi_\Delta(\lambda) = 0$ if $\lambda \notin \Delta$.

Let $\mathcal{B}^{\alpha\beta}_{\theta\gamma}$ be a class of functions F defined on $\sigma_\alpha \times \sigma_\beta$, $\alpha, \beta = 1, 2$, for each $\lambda \in \sigma_\alpha, \mu \in \sigma_\beta$ as operator $F(\lambda, \mu) : \mathcal{G}_\beta(\mu) \to \mathcal{G}_\alpha(\lambda)$, with $\| F \|_\mathcal{B} < \infty$, where

$$\| F \|_\mathcal{B} = \sup_{\substack{\mu \in \sigma_\beta \\ \lambda \in \sigma_\alpha}} (1 + |\lambda|)^\theta (1 + |\mu|)^\theta |F(\lambda, \mu)| +$$

$$+ \sup_{\substack{\lambda, \lambda' \in \sigma^c_\alpha \\ \mu \in \sigma_\beta}} \left\{ (1 + |\mu|)^\theta \frac{|F(\lambda, \mu) - F(\lambda', \mu)|}{|\lambda - \lambda'|^\gamma} \right\} +$$

$$+ \sup_{\substack{\lambda \in \sigma_\alpha \\ \mu, \mu' \in \sigma^c_\beta}} \left\{ (1 + |\lambda|)^\theta \frac{|F(\lambda, \mu) - F(\lambda, \mu')|}{|\mu - \mu'|^\gamma} \right\} +$$

$$+ \sup_{\substack{\lambda, \lambda' \in \sigma^c_\alpha \\ \mu, \mu' \in \sigma^c_\beta}} \left\{ \frac{|F(\lambda, \mu) - F(\lambda', \mu) - F(\lambda, \mu') + F(\lambda', \mu')|}{|\lambda - \lambda'|^\gamma |\mu - \mu'|^\gamma} \right\}.$$

With the norm $\| \cdot \|_\mathcal{B}$ this class will constitute a Banach space. We introduce also the Banach space $\mathcal{M}_{\theta\gamma}(\sigma_\alpha)$ of functions f defined on σ_α with the norm

$$\| f \|_\mathcal{M} = \sup_{\lambda \in \sigma_\alpha} (1 + |\lambda|)^\theta |f(\lambda)| + \sup_{\lambda, \lambda' \in \sigma^c_\alpha} \frac{|f(\lambda) - f(\lambda')|}{|\lambda - \lambda'|^\gamma} < \infty.$$

The value $f(\lambda)$ of the function $f \in \mathcal{M}_{\theta\gamma}(\sigma_\alpha)$ is an operator in $\mathcal{G}_\alpha(\lambda)$.

Let $B_{\alpha\beta}$ be an integral operator with a kernel $B_{\alpha\beta}(\lambda, \mu)$ from the space $\mathcal{B}^{\alpha\beta}_{\theta\gamma}$, $\theta > \frac{1}{2}$, $\frac{1}{2} < \gamma < 1$. We assume that $B_{\alpha\beta}(\lambda, \mu)$ is a compact operator, $B_{\alpha\beta}(\lambda, \mu) : \mathcal{G}_\beta(\mu) \to \mathcal{G}_\alpha(\lambda)$, for each $\lambda \in \sigma_\alpha, \mu \in \sigma_\beta$ and $B_{\alpha\beta}(\lambda, \mu) = 0$ if λ belongs to the boundary of σ^c_α or μ belongs to the boundary of σ^c_β.

With this presuppositions the Hamiltonian $\mathbf{H}$ may be considered as a two-channel variant of the Friedrichs model [20], [21]. Investigation of $\mathbf{H}$ repeats almost literally the analysis from Ref. [21]. Therefore we describe here only final results which are quite analogous to [20], [21]. These results are following.

The operator $\mathbf{H}$ is self-adjoint on the set $\mathcal{D}(\mathbf{H}) = \mathcal{D}(A_1) \oplus \mathcal{D}(A_2)$. Continuous spectrum of $\mathbf{H}$ is situated on the set $\sigma_c(\mathbf{H}) = \sigma^c_1 \cup \sigma^c_2$. Let $\mathbf{H}^c$ be the part of $\mathbf{H}$ acting in the invariant subspace corresponding to continuous spectrum. The operator $\mathbf{H}^c$ is unitary equivalent to the operator $\mathbf{H}_0 = A^{(0)}_1 \oplus A^{(0)}_2$ with $A^{(0)}_\alpha$, $\alpha = 1, 2$, the restriction of the operator A_α on $\mathcal{H}^c_\alpha$. Namely, there exist wave operators $U^{(+)}$ and $U^{(-)}$, $U^{(\pm)} = \begin{pmatrix} u^{(\pm)}_{11} & u^{(\pm)}_{12} \\ u^{(\pm)}_{21} & u^{(\pm)}_{22} \end{pmatrix} = s - \lim_{t \to \mp\infty} \mathrm{e}^{i\mathbf{H}t} \mathrm{e}^{-i\mathbf{H}_0 t}$, with the following properties: $\mathbf{H}U^{(\pm)} = U^{(\pm)}\mathbf{H}_0$, $U^{(\pm)*}U^{(\pm)} = I$, $U^{(\pm)}U^{(\pm)*} = I - P$. Here, P is an orthogonal projector on subspace corresponding to the discrete spectrum $\sigma_d(\mathbf{H})$ of the operator $\mathbf{H}$.

The kernel $u^{(\pm)}_{\alpha\alpha}(\lambda, \lambda')$ of the operator $u^{(\pm)}_{\alpha\alpha}, \alpha = 1, 2$, represents an eigenfunction of the continuous spectrum of the problem (1.2) for $z = \lambda' \pm i0$, $\lambda' \in \sigma^c_\alpha$, and satisfies the integral equation

$$u^{(\pm)}_{\alpha\alpha}(\lambda, \lambda') = I^c_\alpha \delta(\lambda - \lambda') - [(A_\alpha - \lambda' \mp i0)^{-1} V_\alpha(\lambda' \pm i0) u^{(\pm)}_{\alpha\alpha}](\lambda, \lambda'), \tag{2.5}$$

where I_α^c is identity operator in $\mathcal{G}_\alpha^c, \lambda \in \sigma_\alpha$. For each concrete sign (plus or minus) and for each $\lambda' \in \sigma_\alpha^c, \lambda' \notin \sigma_d(\mathbf{H})$ the function $u_{\alpha\alpha}^{(\pm)}(\lambda,\lambda')$ is an unique solution of eq.(2.5) in the class of the distributions

$$f_\alpha^{(\pm)}(\lambda) = I_\alpha \delta(\lambda - \lambda') + \frac{f(\lambda)}{\lambda - \lambda' \mp i0}, \quad f \in \mathcal{M}_{\theta'\gamma'}, \tag{2.6}$$

where $\frac{1}{2} < \theta' < \theta$, $\frac{1}{2} < \gamma' < \gamma$. At the same time

$$u_{\alpha\beta}^{(\pm)}(\lambda,\lambda') = -[(A_\alpha - \lambda' \mp i0)^{-1} B_{\alpha\beta} u_{\beta\beta}^{(\pm)}](\lambda,\lambda'), \quad \beta \neq \alpha,$$

is the problem (1.2) eigenfunction corresponding to $\lambda' \in \sigma_\beta^c$.

The functions $u_{\beta\alpha}^{(\pm)}$, $\alpha,\beta = 1,2$, can be explicitly expressed in terms of kernels of the operator

$$T(z) = B - B(\mathbf{H} - z)^{-1} B, \quad B = \begin{bmatrix} 0 & B_{12} \\ B_{21} & 0 \end{bmatrix}.$$

Corresponding formulae read as

$$u_{\beta\alpha}^{(\pm)}(\lambda,\lambda') = \delta_{\beta\alpha} I_\alpha^c \delta(\lambda - \lambda') - \frac{T_{\beta\alpha}(\mu,\lambda',\lambda' \pm i0)}{\mu - \lambda' \mp i0}, \quad \mu \in \sigma_\beta,\ \lambda' \in \sigma_\alpha^c,$$

with t-matrices

$$T_{\alpha\alpha} = B_{\alpha\beta} \left[z - A_\beta + B_{\beta\alpha}(A_\alpha - z)^{-1} B_{\alpha\beta} \right]^{-1} B_{\beta\alpha}$$

and

$$T_{\beta\alpha} = B_{\beta\alpha} \left[z - A_\alpha + B_{\alpha\beta}(A_\beta - z)^{-1} B_{\beta\alpha} \right]^{-1} (z - A_\alpha) =$$

$$= (z - A_\beta) \left[z - A_\beta + B_{\beta\alpha}(A_\alpha - z)^{-1} B_{\alpha\beta} \right]^{-1} B_{\beta\alpha}, \quad \beta \neq \alpha.$$

Considering the equation for $T(z)$, $T(z) = B - B(A-z)^{-1} T(z)$, $A = A_1 \oplus A_2$, one shows in the same way as in [20], [21] that for all $z \in \mathbf{C} \setminus \sigma(\mathbf{H})$, each kernel $T_{\beta\alpha}(\mu,\lambda,z)$, $\alpha,\beta = 1,2$, belongs to the class $\mathcal{B}_{\theta'\gamma'}^{\beta\alpha}$ with arbitrary θ', γ' such that $\frac{1}{2} < \theta' < \theta$, $\frac{1}{2} < \gamma' < \gamma$. In respect with variable z, the kernel of $T_{\beta\alpha}(z)$ is continuous in the $\mathcal{B}_{\theta'\gamma'}^{\beta\alpha}$–norm right up to the upper and lower borders of the set $\sigma_c(\mathbf{H}) \setminus \sigma_d(\mathbf{H})$.

Scattering operator $S = U^{(-)*} U^{(+)}$ for a system described by the Hamiltonian $\mathbf{H}$ is unitary in $\mathcal{H}_\alpha^c$. It's kernels $s_{\beta\alpha}(\mu,\lambda)$, $\alpha,\beta = 1,2$, are given by expressions

$$s_{\beta\alpha}(\mu,\lambda) = \delta(\mu - \lambda) \left[\delta_{\beta\alpha} I_\alpha^c - 2\pi i\, T_{\beta\alpha}(\mu,\lambda,\lambda + i0) \right]. \tag{2.7}$$

By U_j, $j = 1,2,\ldots$, we denote eigenvectors, $U_j = \{u_1^{(j)}, u_2^{(j)}\}$, $U_j \in \mathcal{D}(\mathbf{H})$, $\| U_j \| = 1$, and by z_j, $z_j \in \mathbf{R}$, the respective eigenvalues of the operator $\mathbf{H}$ discrete spectrum $\sigma_d(\mathbf{H})$. The component $u_\alpha^{(j)}$, $\alpha = 1,2$, of the vector U_j is a solution of Eq. (1.2) at $z = z_j$. If $z_j \in \sigma_\beta^c$ then $(B_{\beta\alpha} u_\alpha^{(j)})(z_j) = 0$.

3. BASIC EQUATION

The paper is devoted to construction of such operator H_α that it's each eigenfunction $u_\alpha, H_\alpha u_\alpha = z u_\alpha$, together with eigenvalue z, satisfies Eq. (1.2). This operator will be

found as a solution of the non-linear operator equation (1.4). To obtain this equation we need the following operator-value function $V_\alpha(Y)$ of the operator variable Y :

$$V_\alpha(Y) = B_{\alpha\beta} \int\limits_{\sigma_\beta} E_\beta(d\mu) B_{\beta\alpha}(Y-\mu)^{-1},$$

$Y : \mathcal{H}_\alpha \to \mathcal{H}_\alpha$. We suppose here that $(Y - \mu I)^{-1} \in L_\infty(\sigma_\beta, \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha))$ if $\mu \in \sigma_\beta$. This means that σ_β has not to be included into the spectrum of the operator Y. Integral $Q(T) = \int\limits_{\sigma_\beta} E_\beta(d\mu) B_{\beta\alpha} T(\mu)$ for $T \in L_\infty(\sigma_\beta, \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha))$, $\|T\|_\infty = E_\beta\!-\sup\limits_{\mu\in\sigma_\beta} \|T(\mu)\| < \infty$, is constructed in the same way as integrals of scalar functions over spectral measure (see Ref. [25], p.130). Namely as a limit value, in respect to the operator norm in $\mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha)$, of respective finite integral sums for piecewise-constant operator-value functions approximating T in $L_\infty(\sigma_\beta, \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha))$. We show the existence of this integral at least in the case when the Hilbert-Schmidt norm $\|B_{\alpha\beta}\|_2$ is finite.

Lemma 1**:** *Let $T \in L_\infty(\sigma_\beta, \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha))$ and $\| B_{\alpha\beta} \|_2 < \infty$. Then the integral $Q(T)$ exists being a bounded operator, $Q(T) : \mathcal{H}_\alpha \to \mathcal{H}_\beta$, $\| Q(T) \| \le \| T \|_\infty \cdot \| B_{\beta\alpha} \|_2$.*

Proof. We prove the Lemma in the diagonal representation (2.2), (2.3). By (2.4) we have

$$(Qf)(\mu) = \mathop{\int\!\!\!\!\!\!\sum}\limits_{\sigma_\alpha} B_{\beta\alpha}(\mu,\lambda)(T(\mu)f)(\lambda)d\lambda$$

for any $f \in \mathcal{H}_\alpha$. It means that

$$|(Qf)(\mu)|^2 \le \mathop{\int\!\!\!\!\!\!\sum}\limits_{\lambda\in\sigma_\alpha} d\lambda \cdot |B_{\beta\alpha}(\mu,\lambda)|^2 \mathop{\int\!\!\!\!\!\!\sum}\limits_{\lambda\in\sigma_\alpha} d\lambda |(T(\mu)f)(\lambda)|^2 =$$

$$= \mathop{\int\!\!\!\!\!\!\sum}\limits_{\lambda\in\sigma_\alpha} d\lambda |B_{\beta\alpha}(\mu,\lambda)|^2 \cdot \|T(\mu)f\|^2 \le \mathop{\int\!\!\!\!\!\!\sum}\limits_{\lambda\in\sigma_\alpha} d\lambda |B_{\beta\alpha}(\mu,\lambda)|^2 \cdot \|T(\mu)\|^2 \cdot \|f\|^2.$$

Hence, integrating over $\mu \in \sigma_\beta$ we come to the relation

$$\| Qf \|^2 \le \| B_{\beta\alpha} \|_2^2 \cdot \| T \|_\infty^2 \cdot \| f \|^2$$

which completes the proof.

Let us suppose that $(H_\alpha - \mu I)^{-1} \in L_\infty(\sigma_\beta, \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\alpha))$. We note that if $H_\alpha\psi_\alpha = z\psi_\alpha$, then automatically

$$V_\alpha(H_\alpha)\psi_\alpha = B_{\alpha\beta} \int\limits_{\sigma_\beta} E_\beta(d\mu) B_{\beta\alpha}(z-\mu)^{-1}\psi_\alpha =$$

$$= B_{\alpha\beta}(z - A_\beta)^{-1} B_{\beta\alpha}\psi_\alpha = V_\alpha(z)\psi_\alpha. \tag{3.1}$$

It follows from (3.1) that H_α satisfies the relation $H_\alpha\psi_\alpha = (A_\alpha + V_\alpha(H_\alpha))\psi_\alpha$ and we can spread this relation over all the linear combinations of H_α eigenfunctions. Supposing that the eigenfunctions system of H_α is dense in $\mathcal{H}_\alpha$ we spread this equation over $\mathcal{D}(A_\alpha)$. As a result we come to the desired *basic equation* (1.4) for H_α (see also Refs. [9], [17]—[19]). Eq. (1.4) means that the construction of the operator H_α comes to the searching for the operator

$$Q_{\beta\alpha} = \int\limits_{\sigma_\beta} E_\beta(d\mu) B_{\beta\alpha}(H_\alpha - \mu)^{-1}. \tag{3.2}$$

Since $H_\alpha = A_\alpha + B_{\alpha\beta}Q_{\beta\alpha}$, we have

$$Q_{\beta\alpha} = \int\limits_{\sigma_\beta} E_\beta(d\mu)B_{\beta\alpha}(A_\alpha + B_{\alpha\beta}Q_{\beta\alpha} - \mu)^{-1}, \quad \beta \neq \alpha. \tag{3.3}$$

In this paper we restrict ourselves to the study of Eq. (3.3) solvability only in the case when spectra σ_1 and σ_2 are separated,

$$d_0 = \mathrm{dist}(\sigma_1, \sigma_2) > 0. \tag{3.4}$$

Using the Lemma 1 and the contracting mapping Theorem, we prove the following:

THEOREM 1: *Let $M_{\beta\alpha}(\delta)$ be a set of bounded operators X, $X : \mathcal{H}_\alpha \to \mathcal{H}_\beta$, satisfying the inequality $\|X\| \leq \delta$ with $\delta > 0$. If this δ and the norm $\|B_{\alpha\beta}\|_2$ satisfy the condition $\|B_{\alpha\beta}\|_2 < d_0 \min\{\frac{1}{1+\delta}, \frac{\delta}{1+\delta^2}\}$, then Eq. (3.3) is uniquely solvable in $M_{\alpha\beta}(\delta)$.*

PROOF. Let

$$F(X) = \int\limits_{\sigma_\beta} E_\beta(d\mu)B_{\beta\alpha}(A_\alpha + B_{\alpha\beta}X - \mu)^{-1} \tag{3.5}$$

with X, the operator from $\mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\beta)$.

Firstly, consider conditions when the function F maps the set $M_\alpha(\delta)$ into itself. We suppose here that $B_{\alpha\beta}$ and X are such that

$$\|B_{\alpha\beta}\|_2\|X\| \leq \delta\|B_{\alpha\beta}\|_2 < d_0 \tag{3.6}$$

and consequently, $\|B_{\alpha\beta}X\| \leq d_0$. This means that spectrum of the operator $A_\alpha + B_{\alpha\beta}X$ does not intersect with the set σ_β. Hence, the resolvent $(A_\alpha + B_{\alpha\beta}X - \mu)^{-1}$ exists and is bounded for any $\mu \in \sigma_\beta$. Thus, by Lemma 1 we have

$$\|F(X)\| \leq \|B_{\alpha\beta}\|_2 \cdot E_\beta\!-\!\sup_{\mu\in\sigma_\beta} \|(A_\alpha + B_{\alpha\beta}X - \mu)^{-1}\|.$$

Due to identity

$$(A_\alpha + B_{\alpha\beta}X - \mu)^{-1} = \left(I + (A_\alpha - \mu)^{-1}B_{\alpha\beta}X\right)^{-1}(A_\alpha - \mu)^{-1}$$

and inequality $\|B_{\alpha\beta}\| \leq \|B_{\alpha\beta}\|_2$ we make estimation

$$\|(A_\alpha + B_{\alpha\beta}X - \mu)^{-1}\| \leq \frac{1}{1 - \|(A_\alpha - \mu)^{-1}\|\|B_{\alpha\beta}\|_2\|X\|}\|(A_\alpha - \mu)^{-1}\| \leq$$

$$\leq \frac{1}{1 - \frac{1}{d_0}\|B_{\alpha\beta}\|_2\delta} \cdot \frac{1}{d_0} = \frac{1}{d_0 - \|B_{\alpha\beta}\|_2\delta}. \tag{3.7}$$

Therefore, the set $M_\alpha(\delta)$ will be mapped by F into itself if $\|B_{\alpha\beta}\|_2$ and δ are such that

$$\|B_{\alpha\beta}\|_2 \cdot \frac{1}{d_0 - \|B_{\alpha\beta}\|_2\delta} \leq \delta. \tag{3.8}$$

Secondly, study conditions for the function F to be a contracting mapping. Now, we consider the difference

$$F(X) - F(Y) = \int\limits_{\sigma_\beta} E_\beta(d\mu)B_{\beta\alpha}\left[(A_\alpha + B_{\alpha\beta}X - \mu)^{-1} - (A_\alpha + B_{\alpha\beta}Y - \mu)^{-1}\right] =$$

$$= \int_{\sigma_\beta} E_\beta(d\mu) B_{\beta\alpha}(A_\alpha + B_{\alpha\beta}X - \mu)^{-1} B_{\alpha\beta}(Y - X)(A_\alpha + B_{\alpha\beta}Y - \mu)^{-1}.$$

Again, by Lemma 1, we have

$$\|F(X) - F(Y)\| \leq$$
$$\leq \|B_{\alpha\beta}\|_2^2 \cdot \sup_{\mu\in\sigma_\beta} \|(A_\alpha + B_{\alpha\beta}X - \mu)^{-1}\| \cdot \sup_{\mu\in\sigma_\beta} \|(A_\alpha + B_{\alpha\beta}Y - \mu)^{-1}\| \cdot \|(Y - X)\|.$$

With (3.7) we come to the estimate

$$\|F(X) - F(Y)\| \leq \|B_{\alpha\beta}\|_2^2 \cdot \frac{1}{(d_0 - \|B_{\alpha\beta}\|_2\delta)^2} \cdot \|Y - X\|.$$

The function F becomes a contracting mapping if

$$\frac{\|B_{\alpha\beta}\|_2^2}{(d_0 - \|B_{\alpha\beta}\|_2\delta)^2} < 1. \tag{3.9}$$

Solving system of the inequalities (3.6), (3.8) and (3.9) we find

$$\|B_{\alpha\beta}\|_2 < d_0 \min\left\{\frac{\delta}{1+\delta^2}, \frac{1}{1+\delta}\right\}$$

and this completes the proof of Theorem 1.

Corollary 1: *Equation (3.3) is uniquely solvable in the unit ball $M_\alpha(1) \subset \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\beta)$ for any $B_{\alpha\beta}$ such that*

$$\|B_{\alpha\beta}\|_2 < \frac{1}{2} d_0. \tag{3.10}$$

To prove the inequality (3.10), note that $\max\limits_{\delta\geq 0} \min\left\{\frac{\delta}{1+\delta^2}, \frac{1}{1+\delta}\right\} = \frac{1}{2}$ (at $\delta = 1$). Hence, if (3.10) takes place then the function (3.5) is a contracting mapping of the unit ball $M_\alpha(1)$ into itself.

Remark. In the proofs of Lemma 1 and Theorem 1 we did not use the assumption about finiteness of the numbers n_α of intervals included in continuous spectra σ_α^c of the operators A_α, $\alpha = 1, 2$. Really, these assertions take place in the case of arbitrary spectrum σ_α.

Finiteness at least of one of the numbers n_1 and n_2 will be used at the moment. If n_1 and/or n_2 are finite and

$$\|B_{\alpha\beta}Q_{\beta\alpha}\| < d_0 = \mathrm{dist}\{\sigma_1, \sigma_2\}, \quad \alpha = 1, 2,\ \beta \neq \alpha, \tag{3.11}$$

we can state that

$$\|(A_\alpha + B_{\alpha\beta}Q_{\beta\alpha} - \mu)^{-1}\| \leq \frac{C_{\alpha\beta}}{1+|\mu|},\ \alpha = 1, 2, \quad \text{at any } \mu \in \sigma_\beta,\ \beta \neq \alpha, \tag{3.12}$$

with some $C_{\alpha\beta} > 0$, $C_{\alpha\beta} \sim 1/(d_0 - \|B_{\alpha\beta}Q_{\beta\alpha}\|)$. Of course this estimate is essential only in the case when σ_β is unbounded. It follows immediately from Eq. (3.3) that if n_1 and/or n_2 are finite then $Q_{\beta\alpha}f_\alpha \in \mathcal{D}(H_\beta) = \mathcal{D}(A_\beta)$ for any $f_\alpha \in \mathcal{H}_\alpha$.

In this case we can rewrite the equation (3.3) in symmetric form

$$Q_{\beta\alpha}A_\alpha - A_\beta Q_{\beta\alpha} + Q_{\beta\alpha}B_{\alpha\beta}Q_{\beta\alpha} = B_{\beta\alpha}. \tag{3.13}$$

To make this, it is sufficient to calculate the expression $Q_{\beta\alpha}H_\alpha - A_\beta Q_{\beta\alpha}$ for both parts of eq.(3.3) having in mind that we apply it to $f_\alpha \in \mathcal{D}(H_\alpha)$. Did, we have

$$Q_{\beta\alpha}H_\alpha - A_\beta Q_{\beta\alpha} = Q_{\beta\alpha}(A_\alpha + B_{\alpha\beta}Q_{\beta\alpha}) - A_\beta Q_{\beta\alpha} = Q_{\beta\alpha}A_\alpha - A_\beta Q_{\beta\alpha} + Q_{\beta\alpha}B_{\alpha\beta}Q_{\beta\alpha}.$$

On the other hand,

$$Q_{\beta\alpha}H_\alpha - A_\beta Q_{\beta\alpha} = \int\limits_{\sigma_\beta} [E_\beta(d\mu)B_{\beta\alpha}(H_\alpha - \mu)^{-1}H_\alpha - \mu E_\beta(d\mu)B_{\beta\alpha}(H_\alpha - \mu)^{-1}] = B_{\beta\alpha}.$$

One finds immediately from Eqs. (3.13), $\alpha = 1, 2$, that if $Q_{\beta\alpha}$ gives solution $H_\alpha = A_\alpha + B_{\alpha\beta}Q_{\beta\alpha}$ of the problem (1.4) in the channel α then

$$Q_{\alpha\beta} = -Q^*_{\beta\alpha} = -\int\limits_{\sigma_\alpha} (H^*_\alpha - \mu)^{-1}B_{\alpha\beta}E_\beta(d\mu) \tag{3.14}$$

gives analogous solution $H_\alpha = A_\alpha + B_{\alpha\beta}Q_{\beta\alpha}$ in the channel β.

THEOREM 2: *Let $Q_{\beta\alpha}$, $Q_{\beta\alpha} \in \mathbf{B}(\mathcal{H}_\alpha, \mathcal{H}_\beta)$, be a solution of Eq. (3.13) satisfying together with $Q_{\alpha\beta} = -Q^*_{\beta\alpha}$ the conditions (3.12). Then the transform $\mathbf{H}' = \mathcal{Q}^{-1}\mathbf{H}\mathcal{Q}$ with $\mathcal{Q} = \begin{bmatrix} I_1 & Q_{12} \\ Q_{21} & I_2 \end{bmatrix}$ reduces the operator $\mathbf{H}$ to the block–diagonal form, $\mathbf{H}' = \mathrm{diag}\{H_1, H_2\}$ where $H_\alpha = A_\alpha + B_{\alpha\beta}Q_{\beta\alpha}$, $\alpha, \beta = 1, 2$, $\beta \neq \alpha$. At the same time, the operators $\mathcal{O}_\alpha = \begin{bmatrix} I_\alpha & 0 \\ Q_{\beta\alpha} & I_\beta \end{bmatrix}$ reduce the Hamiltonian $\mathbf{H}$, $\mathbf{H} = \begin{bmatrix} A_\alpha & B_{\alpha\beta} \\ B_{\beta\alpha} & A_\beta \end{bmatrix}$, to triangular form, $\mathbf{H}^{(\alpha)} \equiv \mathcal{O}_\alpha^{-1}\mathbf{H}\mathcal{O}_\alpha = \begin{bmatrix} H_\alpha & B_{\alpha\beta} \\ 0 & H^*_\beta \end{bmatrix}$.*

PROOFS of both statements are done by direct substituting $\mathcal{Q}$ and $\mathcal{O}_\alpha$ into the definitions of $\mathbf{H}'$ and $\mathbf{H}^{(\alpha)}$ and using the equations (3.13) .

We have to note only that operator $\mathcal{Q}$ is reversible since, according to (3.14),

$$X_\alpha = I_\alpha - Q_{\alpha\beta}Q_{\beta\alpha} = I_\alpha + Q_{\alpha\beta}Q^*_{\alpha\beta} \geq I_\alpha, \quad \alpha = 1, 2, \tag{3.15}$$

and

$$\mathcal{Q}^{-1} = \begin{bmatrix} X_1^{-1} & 0 \\ 0 & X_2^{-1} \end{bmatrix} \cdot \begin{bmatrix} I_1 & -Q_{12} \\ -Q_{21} & I_2 \end{bmatrix}. \tag{3.16}$$

COROLLARY 2: *Subspaces $\mathcal{H}^{(\alpha)} = \mathcal{O}_\alpha(\mathcal{H}_\alpha \oplus \{\mathbf{0}\}) = \{f : f = \{f_\alpha, f_\beta\} \in \mathcal{H}, f_\alpha \in \mathcal{H}_\alpha, f_\beta = Q_{\beta\alpha}f_\alpha\}$ are orthogonal, $\mathcal{H}^{(1)} \perp \mathcal{H}^{(2)}$, and reducing for $\mathbf{H}$, $\mathbf{H}\left(\mathcal{D}(\mathbf{H}) \bigcap \mathcal{H}^{(\alpha)}\right) \subseteq \mathcal{H}^{(\alpha)}$.*

Really, if $f \in \mathcal{H}^{(\alpha)}$, $g \in \mathcal{H}^{(\beta)}$ and $f = \{f_\alpha, Q_{\beta\alpha}f_\alpha\}$, $g = \{Q_{\alpha\beta}g_\beta, g_\beta\}$, then $\langle f, g\rangle = \langle f_\alpha, Q_{\alpha\beta}g_\beta\rangle + \langle Q_{\beta\alpha}f_\alpha, g_\beta\rangle = 0$ since $Q_{\beta\alpha} = -Q^*_{\alpha\beta}$. The invariance of $\mathcal{H}^{(\alpha)}$, $\alpha = 1, 2$, with respect to $\mathbf{H}$ follows from the equality $\mathbf{H}\mathcal{Q} = \mathcal{Q}\mathbf{H}'$.

Assertions quite analogous to the Theorem 2 and Corollary 2 one can find in Refs. [22], [23]. Solvability (for sufficiently small $\|B_{\alpha\beta}\|$) of the equation (3.13) was proved in [22], [23] by rather different method also in the supposition (3.4).

REMARK. It follows from Theorem 2 that operator $\tilde{\mathcal{Q}} = \mathcal{Q}X^{-1/2}$ with $X = \mathrm{diag}\{X_1, X_2\}$ is unitary. Consequently, the operator $\mathbf{H}'' = \tilde{\mathcal{Q}}^*\mathbf{H}\tilde{\mathcal{Q}} = X^{1/2}\mathbf{H}'X^{-1/2}$ becomes self-adjoint in $\mathcal{H}$. Since $\mathbf{H}'' = \mathrm{diag}\{H_1'', H_1''\}$ with $H_\alpha'' = X_\alpha^{1/2}H_\alpha X_\alpha^{-1/2}$, the operators H_α'', $\alpha = 1, 2$, are self-adjoint on $\mathcal{D}(A_\alpha)$ in $\mathcal{H}_\alpha$. Moreover the operators $\mathbf{H}^{(\alpha)} = \tilde{\mathcal{Q}} \cdot \mathrm{diag}\{H_\alpha'', 0\} \cdot \tilde{\mathcal{Q}}^* = \mathcal{Q} \cdot \mathrm{diag}\{H_\alpha, 0\} \cdot \mathcal{Q}^{-1}$ represent parts of the Hamiltonian $\mathbf{H}$ in the corresponding invariant subspaces $\mathcal{H}^{(1)}$ and $\mathcal{H}^{(2)}$ (see also Refs. [22], [23]).

Unfortunately, eigenvectors ψ_α'' of the operators H_α'' differ from those for the initial spectral problem (1.2): ${\psi''}_\alpha = X_\alpha^{1/2}\psi_\alpha$.

LEMMA 2: *Let the kernel $B_{\beta\alpha}(\mu, \lambda)$, $\beta \neq \alpha$, of the operator $B_{\beta\alpha}$ belong to the class $\mathcal{B}_{\theta\gamma}^{\beta\alpha}$ with $\theta > \frac{1}{2}$ and $Q_{\beta\alpha}$ be a solution of Eq. (3.3) satisfying together with $Q_{\alpha\beta} = -Q_{\beta\alpha}^*$ the conditions (3.12). Then*

(a) the operator $Q_{\beta\alpha}$ is an integral operator, $Q_{\beta\alpha} : \mathcal{H}_\alpha \to \mathcal{H}_\beta$, with a kernel $Q_{\beta\alpha}(\mu, \lambda)$ belonging to $\mathcal{B}_{\theta\gamma}^{\beta\alpha}$;

(b) the potential $W_\alpha \equiv B_{\alpha\beta}Q_{\beta\alpha}$ is an integral operator, $W_\alpha : \mathcal{H}_\alpha \to \mathcal{H}_\alpha$, with a kernel $W_\alpha(\lambda, \lambda')$ belonging to $\mathcal{B}_{\theta\gamma}^{\alpha\alpha}$.

PROOF. At the beginning we prove the assertion (b). According to (3.14),

$$W_\alpha = -B_{\alpha\beta} \int_{\sigma_\alpha} (H_\beta^* - \lambda)^{-1} B_{\beta\alpha} E_\alpha(d\lambda) \tag{3.17}$$

with $H_{\beta\alpha}^* = A_\beta + W_\beta^* = A_\beta + Q_{\alpha\beta}^* B_{\alpha\beta}$. Since the inequalities (3.12) take place we write

$$\|(H_\beta^* - \lambda)^{-1}\| = \|(H_\beta - \lambda)^{-1}\| \leq C_{\beta\alpha}$$

for any $\lambda \in \sigma_\alpha$. In the diagonal representation (2.2),(2.3), the equation (3.17) turns in

$$W_\alpha(\lambda, \lambda') = -B_{\alpha\beta}(\lambda, \cdot\,)(H_\beta^* - \lambda)^{-1} B_{\beta\alpha}(\,\cdot\,, \lambda').$$

It means that

$$\begin{aligned} |W_\alpha(\lambda, \lambda')| &\leq \|B_{\alpha\beta}(\lambda, \cdot\,)\|_{\mathcal{H}_\beta} \cdot \|(H_\beta^* - \lambda)^{-1}\| \cdot \|B_{\beta\alpha}(\,\cdot\,, \lambda')\|_{\mathcal{H}_\beta} \leq \\ &\leq C_{\beta\alpha}\|B_{\alpha\beta}(\lambda, \cdot\,)\|_{\mathcal{H}_\beta} \cdot \|B_{\beta\alpha}(\,\cdot\,, \lambda')\|_{\mathcal{H}_\beta}. \end{aligned} \tag{3.18}$$

Here, $\|B_{\alpha\beta}(\lambda, \cdot\,)\|_{\mathcal{H}_\beta} = \left[\oint\limits_{\sigma_\beta} |B_{\alpha\beta}(\lambda, \mu)|^2 d\mu\right]^{1/2}$. Since $\theta > \frac{1}{2}$, we have $\|B_{\alpha\beta}(\lambda, \cdot\,)\|_{\mathcal{H}_\beta} \leq \dfrac{c(\theta)}{(1 + |\lambda|)^\theta} \cdot \|B\|_{\mathcal{B}}$ with some $c(\theta)$, $c(\theta) > 0$, depending only on θ. Analogously,

$$\|B_{\beta\alpha}(\,\cdot\,, \lambda')\|_{\mathcal{H}_\beta} = \|\overline{B_{\alpha\beta}}(\lambda', \cdot\,)\|_{\mathcal{H}_\beta} \leq \frac{c(\theta)}{(1 + |\lambda|)^\theta} \cdot \|B\|_{\mathcal{B}},$$

where the operator $\overline{B_{\alpha\beta}}(\lambda, \mu)$, $\overline{B_{\alpha\beta}}(\lambda, \mu) : \mathcal{G}_\beta(\mu) \to \mathcal{G}_\alpha(\lambda)$, is adjoint to $B_{\beta\alpha}(\mu, \lambda)$.

Estimations similar to (3.18) may be done also for $|W_\alpha(\lambda'', \lambda') - W_\alpha(\lambda, \lambda')|$, $\lambda, \lambda'' \in \sigma_\alpha^c$, $\lambda \in \sigma_\alpha$, $|W_\alpha(\lambda, \lambda''') - W_\alpha(\lambda, \lambda')|$, $\lambda \in \sigma_\alpha$, $\lambda''', \lambda' \in \sigma_\alpha^c$, and $|W_\alpha(\lambda, \lambda') - W_\alpha(\lambda'', \lambda') - W_\alpha(\lambda, \lambda''') + W_\alpha(\lambda'', \lambda''')|$, $\lambda, \lambda', \lambda'', \lambda''' \in \sigma_\alpha^c$, in terms of the norms $\|B_{\alpha\beta}(\lambda, \cdot\,) - B_{\alpha\beta}(\lambda'', \cdot\,)\|_{\mathcal{H}_\beta}$ and $\|B_{\beta\alpha}(\,\cdot\,, \lambda''') - B_{\beta\alpha}(\,\cdot\,, \lambda')\|_{\mathcal{H}_\beta}$. Estimating the latter through $\|B_{\alpha\beta}\|_{\mathcal{B}}$ we come to the inequality

$$\|W_\alpha\|_{\mathcal{B}_{\theta\gamma}^{\alpha\alpha}} \leq c(\theta)C_{\beta\alpha} \cdot \|B_{\beta\alpha}\|^2_{\mathcal{B}_{\theta\gamma}^{\beta\alpha}}$$

with $0 < c(\theta) < \infty$. Therefore, we have proved the assertion (b).

To prove the statement (a) we note that according to (3.3),

$$Q_{\beta\alpha} = \int\limits_{\sigma_\beta} E_\beta(d\mu) B_{\beta\alpha} \left[(A_\alpha - \mu)^{-1} - (H_\alpha - \mu)^{-1} W_\alpha (A_\alpha - \mu)^{-1} \right]$$

or, in the diagonal representation (2.2),(2.3),

$$Q_{\beta\alpha}(\mu, \lambda) = \frac{B_{\beta\alpha}(\mu, \lambda)}{\lambda - \mu} - \frac{B_{\beta\alpha}(\mu, \cdot\,)(H_\alpha - \mu)^{-1} W_\alpha(\,\cdot\,, \lambda)}{\lambda - \mu}.$$

Repeating literally the last part of the proof of the assertion (b) we come to the inequality

$$\|Q_{\beta\alpha}\|_{\mathcal{B}^{\beta\alpha}_{\theta\gamma}} \leq \sup_{\substack{\mu \in \sigma_\beta \\ \lambda \in \sigma_\alpha}} \frac{1}{|\lambda - \mu|} \cdot \Bigg\{ \|B_{\beta\alpha}\|_{\mathcal{B}^{\beta\alpha}_{\theta\gamma}} +$$

$$+ c(\theta) \cdot \|B_{\beta\alpha}\|_{\mathcal{B}^{\beta\alpha}_{\theta\gamma}} \cdot \sup_{\mu \in \sigma_\beta} \|(H_\alpha - \mu)^{-1}\| \cdot \|W_\alpha\|_{\mathcal{B}^{\alpha\alpha}_{\theta\gamma}} \Bigg\}, \quad 0 < c(\theta) < +\infty.$$

Consequently $Q_{\beta\alpha} \in \mathcal{B}^{\beta\alpha}_{\theta\gamma}$ and

$$\|Q_{\beta\alpha}\|_{\mathcal{B}} \leq \frac{1}{d_0} \cdot \left\{ \|B_{\beta\alpha}\|_{\mathcal{B}} + c(\theta) C_{\alpha\beta} C_{\beta\alpha} \cdot \|B_{\beta\alpha}\|^3_{\mathcal{B}} \right\}, \quad 0 < c(\theta) < +\infty.$$

This completes the proof of Lemma 2.

Corollary 3: *If $B_{\beta\alpha} \in \mathcal{B}^{\beta\alpha}_{\theta\gamma}$, $\theta > \frac{1}{2}$, then the solution of Eq. (3.3) described by Theorem 1 belongs to the class $\mathcal{B}^{\beta\alpha}_{\theta\gamma}$, too.*

This statement is based on the fact that the mentioned solution satisfies automatically the conditions (3.11) and, hence, the conditions (3.12).

4. EIGENFUNCTIONS AND THE EXPANSION THEOREM

In the preceding section, we have proved the existence (in the unit ball $M_\alpha(1) \subset \mathcal{H}_\alpha$) of a solution $Q_{\beta\alpha}$ of the basic equation (3.3) only in the case when spectra σ_1, σ_2 of the operators A_1, A_2 are separated, $\text{dist}\{\sigma_1, \sigma_2\} = d_0 > 0$, and $\|B_{12}\|_2 = \|B_{21}\|_2 < \frac{d_0}{2}$. May be, however, Eqs. (1.4) and (3.3) have solutions also in other cases. That is why we study the spectral properties of the operator $H_\alpha = A_\alpha + B_{\alpha\beta} Q_{\beta\alpha}$ not supposing that $\|B_{\alpha\beta}\|_2 < \frac{d_0}{2}$ and using more general requirements (3.12) only, with $C_{\alpha\beta}$, some positive numbers, $\alpha, \beta = 1, 2$, $\beta \neq \alpha$. Of course, we assume again that the condition (3.4) takes place. Remember that the requirements (3.12) are sufficient for existence of the operators $V_\alpha(H_\alpha)$. As well, the equations (3.13) and (3.14) take place and the assertions of Theorem 2 and Lemma 2 are valid.

So, let us suppose that $Q_{\beta\alpha}$ and $Q_{\alpha\beta} = -Q^*_{\beta\alpha}$ are solutions of Eqs. (3.3) and (3.13) satisfying the conditions (3.12). It follows from Lemma 1 that $Q_{\beta\alpha} \in \mathbf{B}_{\beta\alpha}(\mathcal{H}_\alpha, \mathcal{H}_\beta)$ as well

as $Q_{\alpha\beta} \in \mathbf{B}_{\alpha\beta}(\mathcal{H}_\beta, \mathcal{H}_\alpha)$. If $B_{\beta\alpha} \in \mathcal{B}_{\theta\gamma}^{\beta\alpha}$, $\theta > \frac{1}{2}$, then, according to Lemma 2, $Q_{\beta\alpha} \in \mathcal{B}_{\theta\gamma}^{\beta\alpha}$ and $Q_{\alpha\beta} \in \mathcal{B}_{\theta\gamma}^{\alpha\beta}$.

By Theorem 2, the operator $\mathbf{H}' = \operatorname{diag}\{H_1, H_2\}$ is connected with the (self-adjoint) operator $\mathbf{H}$ by a similarity transform. Thus, the spectra $\sigma(H_1)$ and $\sigma(H_2)$ of the operators H_α, $\alpha = 1, 2$, are real and $\sigma(H_1) \bigcup \sigma(H_2) = \sigma(\mathbf{H})$. Continuous spectrum $\sigma_c(H_\alpha)$ of the each operator H_α coincides with that of the operator A_α, $\sigma_c(H_\alpha) = \sigma_\alpha^c$, since due to $\|B_{\alpha\beta}\|_2 < +\infty$, the potential $W_\alpha = B_{\alpha\beta}Q_{\beta\alpha}$ is a compact operator. Since $\sigma_1^c \bigcap \sigma_2^c = \emptyset$ we have $\sigma_c(H_1) \bigcap \sigma_c(H_2) = \emptyset$. We show now that the discrete spectra $\sigma_d(H_\alpha)$, $\alpha = 1, 2$, satisfy a similar condition.

Let us suppose that $\sigma_d(H_\alpha) \neq \emptyset$, $z \in \sigma_d(H_\alpha)$ and ψ_α is the corresponding eigenfunction of H_α, $H_\alpha\psi_\alpha = z\psi_\alpha$, $\psi_\alpha \in \mathcal{D}(H_\alpha) = \mathcal{D}(A_\alpha)$. Then, according to construction of H_α, we have $H_\alpha\psi_\alpha = (A_\alpha + V_\alpha(H_\alpha))\psi_\alpha = (A_\alpha + V_\alpha(z))\psi_\alpha = z\psi_\alpha$. Thus if $z \in \sigma_d(H_\alpha)$ then z becomes automatically a point of the discrete spectrum of the initial spectral problem (1.2). At the same time ψ_α becomes it's eigenfunction.

Let us further denote the eigenfunctions of the operator H_α discrete spectrum by $\psi_\alpha^{(j)}$, $\psi_\alpha^{(j)} = u_\alpha^{(j)}$, keeping for them the same numeration as for eigenvectors of U_j, $U_j = \{u_\alpha^{(j)}, u_\beta^{(j)}\}$, of the Hamiltonian $\mathbf{H}$, $\mathbf{H}U_j = z_jU_j$, $z_j \in \sigma_d(\mathbf{H})$. We assume that in the case of multiple discrete eigenvalues, certain z_j may be repeated in this numeration. By $\mathcal{U}^d$ we denote the set $\mathcal{U}^d = \{U_j, j = 1, 2, \ldots\}$ of all the eigenvectors U_j.

Let $\mathcal{U}_\alpha^d$ be such a subset of $\mathcal{U}^d$ that it's elements have the operator H_α eigenvectors $\psi_\alpha^{(j)}$ in the capacity of the channel α components: $\mathcal{U}_\alpha^d = \{U_j : U_j = \{u_1^{(j)}, u_2^{(j)}\},\ u_\alpha^{(j)} = \psi_\alpha^{(j)}\}$. By Theorem 2, we have $\mathcal{U}_1^d \bigcup \mathcal{U}_2^d = \mathcal{U}^d$.

THEOREM 3: *Let $H_\beta = A_\beta + B_{\beta\alpha}Q_{\alpha\beta}$, correspond (for $\|B_{\beta\alpha}\|_2 < +\infty$) to the same solution $Q_{\alpha\beta} = -Q_{\beta\alpha}^*$ of Eqs. (3.3) and (3.13) as $H_\alpha = A_\alpha + B_{\alpha\beta}Q_{\beta\alpha}$, and the conditions (3.12) are valid. Let $z_j \in \sigma_d(H_\alpha)$ and $H_\alpha u_\alpha^{(j)} = z_j u_\alpha^{(j)}$ with $u_\alpha^{(j)}$, the channel α component of the eigenvector $U_j = \{u_\alpha^{(j)}, u_\beta^{(j)}\}$ of the operator $\mathbf{H}$, $\mathbf{H}U_j = z_jU_j$. Then either $z_j \notin \sigma_d(H_\beta)$, $\beta \neq \alpha$, or (if $z_j \in \sigma_d(H_\beta)$) the vector $u_\beta^{(j)}$ is not an eigenvector of H_β.*

COROLLARY 4: $\mathcal{U}_1^d \bigcap \mathcal{U}_2^d = \emptyset$.

Statement of Theorem 3 means that discrete spectrum $\sigma_d(\mathbf{H})$ is distributed between discrete spectra $\sigma_d(H_1)$ and $\sigma_d(H_2)$ in such a way that operators H_1 and H_2 have not "common" eigenvectors $U_j = \{u_1^{(j)}, u_2^{(j)}\}$: simultaneously, component $u_1^{(j)}$ can not be eigenvector for H_1, and $u_2^{(j)}$ with the same j, for H_2.

PROOF of the Theorem will be given by contradiction.

Let us suppose that $\psi_\alpha^{(j)} = u_\alpha^{(j)}$ is an eigenvector of H_α corresponding to z_j i.e.

$$(A_\alpha + B_{\alpha\beta}Q_{\beta\alpha} - z_j)\psi_\alpha^{(j)} = 0. \tag{4.1}$$

If $z_j \in \sigma_\alpha = \sigma(A_\alpha)$ then automatically $z_j \notin \sigma_d(H_\beta)$ since due to conditions (3.12) we have $\sigma(H_\beta) \bigcap \sigma(A_\alpha) = \emptyset$. Thus in the case when $z_j \in \sigma_\alpha$ the assertion of Theorem is valid.

Let $z_j \notin \sigma(A_\alpha)$. In this case we can rewrite Eq. (4.1) in the form

$$\psi_\alpha^{(j)} = -(A_\alpha - z_j)^{-1} B_{\alpha\beta}Q_{\beta\alpha}\psi_\alpha^{(j)}. \tag{4.2}$$

Let $y_\beta^{(j)} = Q_{\beta\alpha}\psi_\alpha^{(j)}$. It follows from (4.2) that

$$y_\beta^{(j)} + Q_{\beta\alpha}(A_\alpha - z_j)^{-1} y_\beta^{(j)} = 0. \tag{4.3}$$

We will show that the vector $y_\beta^{(j)}$ is a solution of the initial spectral problem (1.2) in the channel β at $z = z_j$ and $\tilde{U}_j = \{\psi_\alpha^{(j)}, y_\beta^{(j)}\}$ is an eigenvector of $\mathbf{H}$, $\mathbf{H}\tilde{U}_j = z_j\tilde{U}_j$. To do this, we act on both parts of Eq. (4.3) by $H_\beta^* - z_j$ remembering that, according to (3.14), $Q_{\beta\alpha} = -Q_{\alpha\beta}^* = -\int_{\sigma_\alpha}(H_\beta^* - \lambda)^{-1}B_{\beta\alpha}E_\alpha(d\lambda)$. We obtain

$$(H_\beta^* - z_j)y_\beta^{(j)} + \int_{\sigma_\alpha}(H_\beta^* - z_j)(H_\beta^* - \lambda)^{-1}(z_j - \lambda)^{-1}B_{\beta\alpha}E_\alpha(d\lambda)B_{\alpha\beta}y_\beta^{(j)} = 0.$$

Using the identity $(H - z)(H - \lambda)^{-1}(z - \lambda)^{-1} = (z - \lambda)^{-1} - (H - \lambda)^{-1}$ we find

$$(H_\beta^* - z_j)y_\beta^{(j)} + \int_{\sigma_\alpha}[(z_j - \lambda)^{-1} - (H_\beta^* - z_j)]B_{\beta\alpha}E_\alpha(d\lambda)B_{\alpha\beta}y_\beta^{(j)} = 0$$

or, and it is the same,

$$(H_\beta^* - z_j)y_\beta^{(j)} - B_{\alpha\beta}(A_\alpha - z_j)^{-1}B_{\alpha\beta}y_\beta^{(j)} + Q_{\beta\alpha}B_{\alpha\beta}y_\beta^{(j)} = 0. \tag{4.4}$$

However $H_\beta^* = A_\beta - Q_{\beta\alpha}B_{\alpha\beta}$. Hence the relation (4.4) turns in equation (1.2) for the channel β,

$$[A_\beta - B_{\beta\alpha}(A_\alpha - z_j)^{-1}B_{\alpha\beta} - z_j]y_\beta^{(j)} = 0.$$

So, we have proved that $y_\beta^{(j)}$ is a solution of the initial problem in the channel β and we did deal with an eigenvector $U_j = \{u_\alpha^{(j)}, u_\beta^{(j)}\}$ of the operator $\mathbf{H}$ having the components $u_\alpha^{(j)} = \psi_\alpha^{(j)}$ and $u_\beta^{(j)} = y_\alpha^{(j)}$.

Let us show that $y_\beta^{(j)}$ can not be an eigenvector of H_β corresponding to the eigenvalue z_j. Actually, due to (4.3) we have

$$\mathrm{a} \equiv \langle y_\beta^{(j)} + Q_{\beta\alpha}(A_\alpha - z_j)^{-1}B_{\alpha\beta}y_\beta^{(j)}, y_\beta^{(j)}\rangle = 0.$$

On the other hand

$$\mathrm{a} = \|y_\beta^{(j)}\|^2 + \langle (A_\alpha - z_j)^{-1}B_{\alpha\beta}y_\beta^{(j)}, Q_{\beta\alpha}^* y_\beta^{(j)}\rangle.$$

If $y_\beta^{(j)}$ is an eigenvector of H_β, $H_\beta y_\beta^{(j)} = z_j y_\beta^{(j)}$, then

$$Q_{\beta\alpha}^* y_\beta^{(j)} = -Q_{\alpha\beta}y_\beta^{(j)} = -\int_{\sigma_\alpha} E_\alpha(d\lambda)B_{\alpha\beta}(H_\beta - \lambda)^{-1}y_\beta^{(j)} = (A_\alpha - z_j)^{-1}B_{\alpha\beta}y_\beta^{(j)}.$$

It means that

$$\mathrm{a} = \|y_\beta^{(j)}\|^2 + \|(A_\alpha - z_j)^{-1}B_{\alpha\beta}y_\beta^{(j)}\|^2 \geq \|y_\beta^{(j)}\|^2.$$

Since $\mathrm{a} = 0$ we get $y_\beta^{(j)} = 0$ and, due to (4.2), $\psi_\alpha^{(j)} = 0$. However, by supposition, $\psi_\alpha^{(j)} \neq 0$. Thus, we come to a contradiction and $y_\beta^{(j)}$ can not be an eigenvector of H_β. And so, if

$z_j \in \sigma_d(H_\alpha)$ and $H_\alpha u_\alpha^{(j)} = z_j u_\alpha^{(j)}$ then $u_\beta^{(j)}$ is not an eigenvector of H_β. The proof of Theorem 3 is completed.

Let us pay attention to the continuous spectrum of H_α assuming here that $B_{\alpha\beta} \in \mathcal{B}_{\theta\gamma}^{\beta\alpha}$, $\theta > \frac{1}{2}$, $\gamma > \frac{1}{2}$, and consequently, $Q_{\alpha\beta} \in \mathcal{B}_{\theta\gamma}^{\beta\alpha}$, $\alpha, \beta = 1, 2$, $\beta \neq \alpha$.

Consider at $\lambda' \in \sigma_\alpha^c$ the integral equations

$$\psi_\alpha^{(\pm)}(\lambda, \lambda') = I_\alpha^c \delta(\lambda - \lambda') - [(A_\alpha - \lambda' \mp i0)^{-1} W_\alpha \psi_\alpha^{(\pm)}](\lambda, \lambda'), \ \alpha = 1, 2, \tag{4.5}$$

where as usually $W_\alpha = B_{\alpha\beta} Q_{\beta\alpha}$. Since $W_\alpha \in \mathcal{B}_{\theta\gamma}^{\alpha\alpha}$, the integral operator with the kernel $\dfrac{W_\alpha(\lambda, \lambda')}{\lambda - \lambda' \mp i0}$ is compact in $\mathcal{M}_{\theta'\gamma'}$, $\dfrac{1}{2} < \theta' < \theta$, $0 < \gamma' < \gamma$ (cf. Refs. [20], [21]). If $\lambda' \notin \sigma_d(H_\alpha)$ then Eq. (4.5) for $\psi_\alpha^{(+)}$ as well as for $\psi_\alpha^{(-)}$ is uniquely solvable (see Ref. [21]) in the class of the form (2.6) distributions.

Denote by $\Psi_\alpha^{(\pm)}$, $\Psi_\alpha^{(\pm)} : \mathcal{H}_\alpha^c \to \mathcal{H}_\alpha$, the integral operator with the kernel $\psi_\alpha^{(\pm)}(\lambda, \lambda')$. The operator $\Psi_\alpha^{(\pm)}$ is bounded and $\Psi_\alpha^{(\pm)} \mathcal{D}(A_\alpha^{(0)}) \subseteq \mathcal{D}(H_\alpha)$ [20], [21]. It follows from (4.5) that $\Psi_\alpha^{(\pm)}$ has the property $H_\alpha \Psi_\alpha^{(\pm)} = \Psi_\alpha^{(\pm)} A_\alpha^{(0)}$. Thus, $Q_{\beta\alpha} \Psi_\alpha^{(\pm)}(\cdot, \lambda') = (\lambda' - A_\beta)^{-1} B_{\beta\alpha} \Psi_\alpha^{(\pm)}(\cdot, \lambda')$. Substitution of this expression in (4.5) shows that $\psi_\alpha^{(\pm)}$ satisfies (2.5). Due to the uniqueness of Eq. (2.5) solution at $\lambda' \notin \sigma_d(\mathbf{H})$ we have $\psi_\alpha^{(\pm)}(\lambda, \lambda') = u_{\alpha\alpha}^{(\pm)}(\lambda, \lambda')$. This means that each eigenfunction $u_{\alpha\alpha}^{(\pm)}(\lambda, \lambda')$, $\lambda' \in \sigma_\alpha^c$, $\lambda' \notin \sigma_d(\mathbf{H})$ of the initial spectral problem problem (1.2) is also an eigenfunction of H_α.

Consider the functions $\tilde{\psi}_\alpha^{(j)} = \psi_\alpha^{(j)} - Q_{\alpha\beta} u_\beta^{(j)}$ and $\tilde{\psi}_\alpha^{(\pm)}(\cdot, \lambda') = \psi_\alpha^{(\pm)}(\cdot, \lambda') - Q_{\alpha\beta} u_{\beta\alpha}^{(\pm)}(\cdot, \lambda')$, $\lambda' \in \sigma_\alpha^c$. Let $\tilde{\Psi}_\alpha^{(\pm)}$, $\tilde{\Psi}_\alpha^{(\pm)} : \mathcal{H}_\alpha^c \to \mathcal{H}_\alpha$, be the integral operator with the kernel $\tilde{\psi}_\alpha^{(\pm)}(\lambda, \lambda')$.

Theorem 4: *The functions $\tilde{\psi}_\alpha^{(j)}$ (with j such that $U_j \in \mathcal{U}_\alpha^d$) are eigenfunctions of adjoint operator $H_\alpha^*, H_\alpha^* = A_\alpha + Q_{\beta\alpha}^* B_{\beta\alpha}$, discrete spectrum, $H_\alpha^* \tilde{\psi}_\alpha^{(j)} = z_j \tilde{\psi}_\alpha^{(j)}$. Operators $\tilde{\Psi}_\alpha^{(\pm)}$ have the property $H_\alpha^* \tilde{\Psi}_\alpha^{(\pm)} = \tilde{\Psi}_\alpha^{(\pm)} A_\alpha^{(0)}$. At the same time the orthogonality relations take place: $\langle \psi_\alpha^{(j)}, \tilde{\psi}_\alpha^{(k)} \rangle = \delta_{jk}$, $\Psi_\alpha^{(\pm)*} \tilde{\Psi}_\alpha^{(\pm)} = I_\alpha|_{\mathcal{H}_\alpha^c}$, $\tilde{\Psi}_\alpha^{(\pm)*} \psi_\alpha^{(j)} = 0$ and $\Psi_\alpha^{(\pm)*} \tilde{\psi}_\alpha^{(j)} = 0$. Also, the following completeness relations are valid,*

$$\sum_{j:\, U_j \in \mathcal{U}_\alpha^d} \psi_\alpha^{(j)} \langle \cdot, \psi_\alpha^{(j)} \rangle + \Psi_\alpha^{(\pm)} \tilde{\Psi}_\alpha^{(\pm)*} = I_\alpha, \quad \alpha = 1, 2, \tag{4.6}$$

Proof. Show for example that

$$H_\alpha^* \tilde{\psi}_\alpha^{(j)} = z_j \tilde{\psi}_\alpha^{(j)} \tag{4.7}$$

(remember that $z_j \in \mathbf{R}$). We have

$$H_\alpha^* \tilde{\psi}_\alpha^{(j)} = (A_\alpha - Q_{\alpha\beta} B_{\beta\alpha})(\tilde{\psi}_\alpha^{(j)} - Q_{\alpha\beta} u_\beta^{(j)}) =$$

$$= (A_\alpha - Q_{\alpha\beta} B_{\beta\alpha}) \psi_\alpha^{(j)} - (A_\alpha Q_{\alpha\beta} - Q_{\alpha\beta} B_{\beta\alpha} Q_{\alpha\beta}) u_\beta^{(j)}. \tag{4.8}$$

Note that $A_\alpha = H_\alpha - B_{\alpha\beta} Q_{\beta\alpha}$ and, hence,

$$(A_\alpha - Q_{\alpha\beta} B_{\beta\alpha}) \psi_\alpha^{(j)} = z_j \psi_\alpha^{(j)} - (B_{\alpha\beta} Q_{\beta\alpha} + Q_{\alpha\beta} B_{\beta\alpha}) \psi_\alpha^{(j)}. \tag{4.9}$$

Second term in the right part of (4.9) may be easily expressed through $u_\beta^{(j)}$. Actually, $u_\beta^{(j)} = -(A_\beta - z_j)^{-1} B_{\beta\alpha} \psi_\alpha^{(j)}$ (we use again the property $\sigma(H_\alpha) \bigcap \sigma_\beta = \emptyset$ following

from (3.12)). Since Eqs. (3.2) and $H_\alpha\psi_\alpha^{(j)} = z_j\psi_\alpha^{(j)}$ take place, we find $Q_{\beta\alpha}\psi_\alpha^{(j)} = B_{\alpha\beta}u_\beta^{(j)}$. Consequently,

$$(B_{\alpha\beta}Q_{\beta\alpha} + Q_{\alpha\beta}B_{\beta\alpha})\psi_\alpha^{(j)} = B_{\beta\alpha}Q_{\beta\alpha}\psi_\alpha^{(j)} + Q_{\alpha\beta}(A_\beta - z_j)(A_\beta - z_j)^{-1}B_{\beta\alpha}\psi_\alpha^{(j)} =$$
$$= B_{\beta\alpha}u_\beta^{(j)} - Q_{\alpha\beta}(A_\beta - z_j)u_\beta^{(j)}.$$

Substituting the expressions obtained into (4.9) and then into (4.8), we get

$$H_\alpha^*\tilde{\psi}_\alpha^{(j)} = z_j(\psi_\alpha^{(j)} - Q_{\alpha\beta}u_\beta^{(j)}) + [-B_{\alpha\beta} + Q_{\alpha\beta}A_\beta - A_\alpha Q_{\alpha\beta} + Q_{\alpha\beta}B_{\beta\alpha}Q_{\alpha\beta}]u_\beta^{(j)}.$$

According to the equations (3.13), the expression in the square brackets is equal to zero and we come to (4.7).

The equalities $H_\alpha^*\tilde{\psi}_\alpha^{(\pm)}(\,\cdot\,,\lambda') = \lambda'\tilde{\psi}_\alpha^{(\pm)}(\,\cdot\,,\lambda')$, $\lambda' \in \sigma_\alpha^c$, are proved quite analogously.

The orthogonality relations $\langle\psi_\alpha^{(j)}, \tilde{\psi}_\alpha^{(k)}\rangle = \delta_{jk}$, $\tilde{\Psi}_\alpha^{(\pm)*}\psi_\alpha^{(j)} = 0$ and $\Psi_\alpha^{(\pm)*}\tilde{\psi}_\alpha^{(j)} = 0$ are trivial. Proofs of the relation $\Psi_\alpha^{(\pm)*}\tilde{\Psi}_\alpha^{(\pm)} = I_\alpha|_{\mathcal{H}_\alpha^c}$, and the equality (4.6) are very similar. Both these proofs are based on use of properties of the wave operators $U^{(\pm)}$. As a sample, we give a proof of the completeness relation (4.6).

Consider the operator

$$\mathcal{A} = \sum_{j:\,U_j\in\mathcal{U}_\alpha^d} \psi_\alpha^{(j)}\langle\,\cdot\,,\tilde{\psi}_\alpha^{(j)}\rangle + \Psi_\alpha\tilde{\Psi}_\alpha^* =$$
$$= \sum_{j:\,U_j\in\mathcal{U}_\alpha^d} \psi_\alpha^{(j)}\langle\,\cdot\,,\psi_\alpha^{(j)} - Q_{\alpha\beta}u_\beta^{(j)}\rangle + \Psi_\alpha[\Psi_\alpha^* - (Q_{\alpha\beta}u_{\beta\alpha})^*]. \tag{4.10}$$

For convenience, we omit signs "±" in notations of $\Psi_\alpha^{\pm)} \equiv u_{\alpha\alpha}^{(\pm)}$, $u_{\beta\alpha}^{(\pm)}$ and $\tilde{\Psi}_\alpha^{\pm)}$ taking in mind for example the case of sign "+". We have from (4.10):

$$\mathcal{A} = \sum_{j:\,U_j\in\mathcal{U}_\alpha^d} \psi_\alpha^{(j)}\langle\,\cdot\,,\psi_\alpha^{(j)}\rangle + \Psi_\alpha\Psi_\alpha^* - \sum_{j:\,U_j\in\mathcal{U}_\alpha^d} \psi_\alpha^{(j)}\langle\,\cdot\,,Q_{\alpha\beta}u_\beta^{(j)}\rangle - \Psi_\alpha u_{\beta\alpha}^* Q_{\alpha\beta}^*.$$

It follows from the completeness relations $U^{(\pm)*}U^{(\pm)} = I - P$ for wave operators $U^{(\pm)}$ that

$$\Psi_\alpha u_{\beta\alpha}^* \equiv u_{\alpha\alpha}u_{\beta\alpha}^* = -u_{\alpha\beta}u_{\beta\beta}^* - \sum_{z_j\in\sigma(\mathbf{H})} u_\alpha^{(j)}\langle\,\cdot\,,u_\beta^{(j)}\rangle. \tag{4.11}$$

Since $u_{\beta\beta}^*Q_{\alpha\beta}^* = (Q_{\alpha\beta}u_{\beta\beta})^* = u_{\alpha\beta}^*$, we can write with a help of (4.11) that

$$\mathcal{A} = u_{\alpha\alpha}u_{\alpha\alpha}^* + u_{\alpha\beta}u_{\alpha\beta}^* + \sum_{j:\,U_j\in\mathcal{U}_\alpha^d} \psi_\alpha^{(j)}\langle\,\cdot\,,\psi_\alpha^{(j)}\rangle + \sum_{\substack{z_j\,\in\,\sigma_d(\mathbf{H})\\ U_j\,\notin\,\mathcal{U}_\alpha^d}} u_\alpha^{(j)}\langle\,\cdot\,,Q_{\alpha\beta}u_\beta^{(j)}\rangle.$$

In the last sum, the conditions $z_j \in \sigma_d(\mathbf{H})$ and $U_j \notin \mathcal{U}_\alpha^d$ mean really that we deal with any j such that $U_j \in \mathcal{U}_\beta^d$. This follows from the equalities $\mathcal{U}_1^d \bigcup \mathcal{U}_2^d = \mathcal{U}^d$ and $\mathcal{U}_1^d \bigcap \mathcal{U}_2^d = \emptyset$ (see Theorem 3 and Corollary 4). For $U_j \in \mathcal{U}_\beta^d$, the vector $u_\beta^{(j)}$ is eigenfunction of H_β, $u_\beta^{(j)} = \psi_\beta^{(j)}$, and $Q_{\alpha\beta}u_\beta^{(j)} = Q_{\alpha\beta}\psi_\beta^{(j)} = u_\alpha^{(j)}$. Thus, $\mathcal{A}$ turns in

$$\mathcal{A} = u_{\alpha\alpha}u_{\alpha\alpha}^* + u_{\alpha\beta}u_{\beta\alpha}^* + \sum_{z_j\in\sigma_d(\mathbf{H})} u_\alpha^{(j)}\langle\,\cdot\,,u_\alpha^{(j)}\rangle = (U^{(\pm)}U^{(\pm)*} + P)_{\alpha\alpha}.$$

Since $U^{(\pm)}U^{(\pm)*} + P = I$ we find $\mathcal{A} = I_\alpha$ and this completes the proof of Theorem 4.

Theorem 4 means in particular that *part H_α^c of operator H_α acting in the invariant subspace corresponding to it's continuous spectrum σ_α^c, is similar to the operator $A_\alpha^{(0)}$, $H_\alpha^c = \Psi_\alpha^{(\pm)}A_\alpha^{(0)}\tilde{\Psi}_\alpha^{(\pm)*}$, and spectrum σ_α^c is absolutely continuous.*

5. INNER PRODUCT MAKING NEW HAMILTONIANS SELF-ADJOINT

We introduce now a new inner product $[\,.\,,\,.\,]_\alpha$ in $\mathcal{H}_\alpha$, $[f_\alpha, g_\alpha]_\alpha = \langle X_\alpha f_\alpha, g_\alpha\rangle$, $f_\alpha, g_\alpha \in \mathcal{H}_\alpha$, with X_α defined as in Theorem 2, $X_\alpha = I_\alpha + Q_{\alpha\beta}Q^*_{\alpha\beta}$, $\alpha = 1, 2$. The operator X_α is positive definite, $X_\alpha \geq I_\alpha$. This means that $[\,.\,,\,.\,]_\alpha$ satisfies all the axioms of inner product.

Theorem 5: *The operator H_α, $\alpha = 1, 2$, is self-adjoint on $\mathcal{D}(A_\alpha)$ with respect to the inner product $[\,.\,,\,.\,]_\alpha$.*

Proof. It follows from Theorem 2 that operator $\mathbf{H}'$ is self-adjoint in $\mathcal{H} = \mathcal{H}_1 \oplus \mathcal{H}_2$ with respect to the inner product $[\,.\,,\,.\,]$, $[f, g] = [Xf, g]$ with $X = \mathrm{diag}\{X_1, X_2\}$. Did, since $\mathcal{Q}^{-1} = \mathcal{Q}^* X^{-1} = X^{-1}\mathcal{Q}^*$, we have for $f, g \in \mathcal{D}(\mathbf{H}') = \mathcal{D}(\mathbf{H}) = \mathcal{D}(A_1) \oplus \mathcal{D}(A_2)$:

$$[\mathbf{H}'f, g] = \langle X\mathcal{Q}^{-1}\mathbf{H}\mathcal{Q}f, g\rangle = \langle X \cdot X^{-1}\mathcal{Q}^*\mathbf{H}\mathcal{Q}f, g\rangle =$$

$$= \langle f, \mathcal{Q}^*\mathbf{H}\mathcal{Q}g\rangle = \langle f, X \cdot X^{-1}\mathcal{Q}^*\mathbf{H}\mathcal{Q}g\rangle = [f, \mathbf{H}'g].$$

Here, we used the fact that in the case of (3.12), $\mathcal{Q}f \in \mathcal{D}(A_1) \oplus \mathcal{D}(A_2)$ if $f \in \mathcal{D}(A_1) \oplus \mathcal{D}(A_2)$.

Taking elements f, g in the equality $[\mathbf{H}'f, g] = [f, \mathbf{H}'g]$ in the form $f = \{f_1, 0\}$, $g = \{g_1, 0\}$ or $f = \{0, f_2\}$, $g = \{0, g_2\}$ with one of the components equal to zero and $f_\alpha, g_\alpha \in \mathcal{D}(A_\alpha)$, $\alpha = 1, 2$, one comes to the statement of Theorem.

Remark. This Theorem may be proved also in another way making use of the equality

$$I_\alpha + Q_{\alpha\beta}Q^*_{\alpha\beta} = \sum_{j:\, U_j \in \mathcal{U}^d_\alpha} \tilde{\psi}^{(j)}_\alpha \langle\,\cdot\,, \tilde{\psi}^{(j)}_\alpha\rangle + \tilde{\Psi}^{\pm)}_\alpha \tilde{\Psi}^{\pm)*}_\alpha \tag{5.1}$$

which is valid for both signs "+" and "−". In this case, a self-adjointness of H_α with respect to $[\,\cdot\,,\,\cdot\,]_\alpha$ follows from the fact that it's spectrum is real and also from relations $H^*_\alpha \tilde{\Psi}^{(\pm)}_\alpha \tilde{\Psi}^{(\pm)*}_\alpha = \tilde{\Psi}^{(\pm)}_\alpha A^{(0)}_\alpha \tilde{\Psi}^{(\pm)*}_\alpha = \tilde{\Psi}^{(\pm)}_\alpha \tilde{\Psi}^{(\pm)*}_\alpha H_\alpha$. The equality (5.1) itself is proved by calculating it's right part in the same way as it was done when the completeness relations (4.6) were established (see proof of Theorem 4).

6. SCATTERING PROBLEM

We establish now that operators $\Psi^{(+)}_\alpha$ and $\Psi^{(-)}_\alpha$ play the same important role describing a time asymptotics of solutions of the Schrödinger equation

$$i\frac{d}{dt}f_\alpha(t) = H_\alpha f_\alpha(t) \tag{6.1}$$

as in the usual self-adjoint case [20], [21].

Theorem 6: *Operator $U_\alpha(t) = \exp(iH_\alpha t)\exp(-iA^{(0)}_\alpha t)$ converges strongly if $t \to \mp\infty$, with respect to the norm $\|\cdot\|^X_\alpha$ corresponding to the inner product $[\,\cdot\,,\,\cdot\,]_\alpha$ in $\mathcal{H}_\alpha$. The limit is equal to* $s-\lim\limits_{t\to\mp\infty} U_\alpha(t) = \Psi^{(\pm)}_\alpha$.

Since the norms $\|\cdot\|^X_\alpha$ and $\|\cdot\|$ in $\mathcal{H}_\alpha$ are equivalent, $\|f\| \leq \|f\|^X_\alpha \leq (1 + \|Q_{\alpha\beta}\| \cdot \|Q_{\beta\alpha}\|)^{1/2}\|f\|$, the same statement takes place also with respect to the initial norm $\|\cdot\|$.

Theorem 7: *For any element $f_\alpha^{(-)} \in \mathcal{H}_\alpha^c$ one can find such unique element $f_\alpha^{(0)}$ that solution $f_\alpha(t) = \exp(-iH_\alpha t) f_\alpha^{(0)}$ of Eq. (6.1) satisfies the asymptotic condition*

$$\lim_{t\to-\infty} \| f_\alpha(t) - \exp(-iA_\alpha^{(0)}t) f_\alpha^{(-)} \|_\alpha^X = 0.$$

There exists the unique element $f_\alpha^{(+)} \in \mathcal{H}_\alpha^c$ such that

$$\lim_{t\to+\infty} \| f_\alpha(t) - \exp(-iA_\alpha^{(0)}t) f_\alpha^{(+)} \|_\alpha^X = 0.$$

Elements $f_\alpha^{(-)}$ and $f_\alpha^{(+)}$ are connected by the relation $f_\alpha^{(+)} = S^{(\alpha)} f_\alpha^{(-)}$ with $S^{(\alpha)} = \Psi_\alpha^{(-)-1}\Psi_\alpha^{(+)} = \tilde{\Psi}_\alpha^{(-)*}\Psi_\alpha^{(+)} = \Psi_\alpha^{(-)*} X_\alpha \Psi_\alpha^{(+)}$.

We do not give here proofs of the Theorems 6 and 7 because they are exactly the same as in the case of one-particle Schrödinger operator in Ref. [26]

Theorem 7 gives the non-stationary formulation of the scattering problem for a system described by Hamiltonian H_α. Moreover $S^{(\alpha)}$ is a scattering operator for this system.

Theorem 8: *Scattering operator $S^{(\alpha)}$ coincides with the component $s_{\alpha\alpha}$ of the scattering operator S, $S = U^{(-)*}U^{(+)}$, for a system described by the two–channel Hamiltonian $\mathbf{H}$.*

Proof. Let us show that operator $S^{(\alpha)}$ has the kernel $s_{\alpha\alpha}(\lambda, \lambda')$ given by Eq. (2.7). To do this, remember that $\tilde{\Psi}_\alpha^{(-)} = \Psi_\alpha^{(-)} - Q_{\alpha\beta} u_{\beta\alpha}^{(-)}$ (see Theorem 4). Therefore,

$$S^{(\alpha)} = (\Psi_\alpha^{(-)*} - u_{\beta\alpha}^{(-)*} Q_{\alpha\beta}^*)\Psi_\alpha^{(+)} = \Psi_\alpha^{(-)*}\Psi_\alpha^{(+)} + u_{\beta\alpha}^{(-)*} Q_{\beta\alpha}\Psi_\alpha^{(+)} = u_{\alpha\alpha}^{(-)*}u_{\alpha\alpha}^{(+)} + u_{\beta\alpha}^{(-)*}u_{\beta\alpha}^{(+)}.$$

Here, we have used the properties $\Psi_\alpha^{\pm)} = u_{\alpha\alpha}^{(\pm)}$, $Q_{\alpha\beta}^* = -Q_{\alpha\beta}$ and $Q_{\beta\alpha}\Psi_\alpha^{(+)} = u_{\beta\alpha}^{(+)}$ established above. Since

$$u_{\alpha\alpha}^{(-)*}u_{\alpha\alpha}^{(+)} + u_{\beta\alpha}^{(-)*}u_{\beta\alpha}^{(+)} = \left(U^{(-)*}U^{(+)}\right)_{\alpha\alpha} = s_{\alpha\alpha},$$

we come to the statement of Theorem. The proof of Theorem 7 is completed.

A kernel of the scattering operator $S^{(\alpha)}$ may be presented also in a usual way (2.7) in terms of the t-matrix $t_\alpha(z) = W_\alpha - W_\alpha(H_\alpha - z)^{-1}W_\alpha$, taken on the energy–shell. Note that $t_\alpha(z)$ differs from $T_{\alpha\alpha}(z)$ introduced in Sec. 2. Did, easy calculations show that

$$t_\alpha(z) = B_{\alpha\beta}[I_\beta + Q_{\beta\alpha}(A_\alpha - z)^{-1}B_{\alpha\beta}]^{-1}Q_{\beta\alpha}. \tag{6.2}$$

Using the basic equation (3.13) one can rewrite (6.2) in the form

$$t_\alpha(z) = T_{\alpha\alpha}(z) + \tilde{t}_\alpha(z)$$

where

$$\tilde{t}_\alpha(z) = B_{\alpha\beta}[A_\beta - B_{\beta\alpha}(A_\alpha - z)^{-1}B_{\alpha\beta}]^{-1}Q_{\alpha\beta}(A_\alpha - z) \not\equiv 0.$$

However the additional term $\tilde{t}_\alpha(z)$ is evidently disappearing on the energy-shell due to presence of the difference $A_\alpha - z$ as an end factor. Actually, in the diagonal representation (2.2),(2.3), $A_\alpha - z$ acts as the factor $\lambda - z$ vanishing at $z = \lambda + i0$. Therefore, kernels of t-matrices t_α and $T_{\alpha\alpha}$ coincide on the energy surface.

Note also that in our case $\sigma_1^c \bigcap \sigma_2^c = \emptyset$. Hence we have $s_{\beta\alpha} = 0$ and $S^{(\alpha)} = s_{\alpha\alpha}$ is unitary.

7. ON USE OF THE TWO–BODY ENERGY–DEPENDENT POTENTIALS IN FEW–BODY PROBLEMS

There is a rather conceptual question (see for instance Refs. [7], [10]) concerning a use of the two–body energy–dependent potentials in few–body non–relativistic scattering problems. Evidently this question is strongly related to the subject of the paper and we will discuss here three approaches seemed to be reasonable when one tries to embed energy–dependent potentials in few–body equations.

A customary way to embed such potentials into the center-of-mass frame N–body Schrödinger equation consists in the following. Namely, one replaces (see papers [6], [8], [12], [13] and Refs. therein) the pair energy z_{ij}, argument of the potential $V_{ij}(z_{ij})$, $i \neq j$, describing interaction in two-body subsystem $\{i, j\}$ (i, j stand for numbers of particles, $i, j = 1, 2, ...,$N) with the difference $Z - T'_{ij}$ between total energy Z of system and the kinetic energy operator T'_{ij} of particles, supplementary to the subsystem $\{i, j\}$. For the resolvent-like energy dependent potentials (1.3) this replacement is quite correct from mathematical point of view since one can reconstruct underlying multichannel (for instance, four–channel if N=3) self-adjoint Hamiltonian [12], [13]. Reducing the spectral problem for this Hamiltonian to the external channel only one gets the N–body Schrödinger equation exactly with the pair potentials $V_{ij}(Z - T'_{ij})$. Thus one can guarantee that spectrum of this equation is real and the scattering problem for the N–body system can be based.

However the replacements $z_{ij} \to Z - T'_{ij}$ meet serious conceptual objections formulated in concentrated form by E.W.Schmid [10]. Did, it follows from the energy conservation law that to obtain a share of total energy belonging to subsystem $\{i, j\}$, one has to subtract from Z not only T'_{ij} but also a potential energy of interaction between particles i, j and the rest particles of the system. This idea shows really a first way for the correct (in the context of Ref. [10]) embedding two-body potentials into N-body equations: one has to redefine pair potentials as solutions $\mathbf{V}_{ij}$ of the following system of equations:

$$\mathbf{V}_{ij} = V_{ij}(Z - T'_{ij} - \sum_{\{i',j'\} \neq \{i,j\}} \mathbf{V}_{i'j'}) \tag{7.1}$$

where $\{i, j\}$ runs all the pair subsystems. So, the usual embbedings $V_{ij}(z_{ij}) \to V_{ij}(Z - T'_{ij})$ may be considered only as a zero approximation to solutions $\mathbf{V}_{ij}(Z)$ of the system (7.1). Unfortunately, this system may be treated relatively easy only in the case of linear dependence of the potentials $V_{ij}(z_{ij})$ on the (pair) energies z_{ij}. One can show in this case that operator–value functions $V_{ij}(Y)$ of the operator variable Y, $Y : L_2(\mathbf{R}^{3(N-1)}) \to L_2(\mathbf{R}^{3(N-1)})$ may be defined in such a way that solutions of Eqs. (7.1) generate only real spectrum for the N–body Schrödinger equation.

In the case of the resolvent-like energy dependence (1.3) of pair interactions one meets serious obstacles in solving the system (7.1) connected with a strong non-linearity of it's equations. Also, no underlying self–adjoint Hamiltonian is still found.

Another way to deal with the two-body energy-dependent potentials in few–body problems is to replace them with energy–independent ones. In fact, in the present work we realized namely this idea which was pronounced by B.H.J.McKellar and C.M.McKay [7]. Did, let us denote now a "share" of the total energy of the N-body system belonging to the pair subsystem $\{i, j\}$, by h_{ij}. Then this h_{ij} has to satisfy the operator equation following from the energy conservation law, too,

$$h_{ij} = h_{ij}^{(0)} + v_{ij} + V_{ij}(h_{ij}), \tag{7.2}$$

where $h_{ij}^{(0)}$ stands for the kinetic energy operator of the pair $\{i,j\}$ and v_{ij}, for an energy-independent part of the pair interaction. Remember that the equation (7.2) in notation (1.4) was a main subject of the present work. If solutions h_{ij} of equations (7.2) be known, one could substitute the (energy-independent) operators $W_{ij} \equiv V_{ij}(h^{ij})$ in the N-body Hamiltonian treating them in conventional way as additional energy-independent potentials. It should be noted however that the potentials W_{ij} are not totally equivalent to the potentials $V_{ij}(z)$ given by (1.3) since the Hamiltonian h_{ij} being solution of (7.2), reproduces only a part of spectrum of the Schrödinger equation

$$(h_{ij}^{(0)} + v_{ij} + V_{ij}(z))\psi = z\psi \tag{7.3}$$

(see Sec. 4). Forbidden eigenstates correspond normally to the spectrum generated by respective internal Hamiltonian [17]. There is also another question: is the spectrum of the N–body Hamiltonian real if potentials W_{ij} are substituted in? Thing is that Hamiltonian h_{ij} becomes self-adjoint only with respect to a new inner product in $L_2(\mathbf{R}^3)$ (see Sec. 5). One can overcome this difficulty replacing h_{ij} with similar Hamiltonian $h'_{ij} = X_{ij}^{1/2} h_{ij} X_{ij}^{-1/2}$ where X_{ij} is analog for h_{ij} of the operators X_α introduced in Sec. 3. Writing h'_{ij} in the form $h'_{ij} = h_{ij}^{(0)} + V'_{ij}$ one gets a new pair potential V'_{ij} which is already self–adjoint with respect to the standard inner product in $L_2(\mathbf{R}^3)$. Thus, one may use then the potentials V'_{ij} being sure that the N–body Hamiltonian constructed is Hermitic. Emphasize that potential V'_{ij} gives the same two–body spectrum and phase shifts as the potential $v_{ij} + W_{ij}$ because h'_{ij} is obtained from h_{ij} by similarity transform. It follows from Theorems 7 and 8 that the phase shifts given by V'_{ij} coincide also with those for Eq. (7.3). Therefore, the operator V'_{ij} turns out one of the phase–equivalent potentials for the two–body subsystem concerned.

So, we have discussed three different approaches to embedding the two-body energy–dependent potentials in few–body problems. Certainly, the approaches based on solving the non-linear equations (7.1) and (7.2) do not seem to be too attractive from the computational point of view. However, in the cases when the internal Hamiltonians of pair subsystems have a finite discrete spectrum only and the coupling of channels is relatively small (see Corollary 1 to Theorem 1), the approach based on solving Eqs. (7.2) could be quite realized numerically.

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